Gauge/Gravity Duality, Winter 2024/25

Lecturer: Dr. Falk Hassler, falk.hassler@uwr.edu.pl Tutorials: M.Sc. Luca Scala, luca.scala@uwr.edu.pl



## 4. Einstein's fields equations and AdS space (2)

To be discussed on Friday,  $25^{\text{th}}$  October, 2024 in the tutorial. Please indicate your preferences until Sunday, 20/10/2024, 20:00:00 on the website.

## Exercise 4.1: The energy-momentum tensor revisited

In the lecture, we learned that the Einstein-Hilbert action coupled to matter gives rise to the field equations

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = \kappa^2 T_{\mu\nu} \tag{1}$$

where the energy momentum tensor in curved space is defined by

$$-\frac{2}{\sqrt{-g}}\frac{\delta S_{\text{matter}}}{\delta g^{\mu\nu}}\,.\tag{2}$$

- a) (2 points) Eliminate the curvature scalar R from the field equations (1)by computing and removing its trace.
- b) (3 points) Use the alternative definition of the energy-momentum tensor (2) and determine  $T_{\mu\nu}$  for a real scalar field  $\phi$  with mass m and the interaction

$$L_{\rm int}(\phi) = -\frac{g_n}{n!}\phi^n$$

c) (3 points) Compute the trace of the energy-momentum tensor, if we consider a massless free theory with the additional contribution

$$S_{R\phi^2} = -\xi \int \mathrm{d}^d x \sqrt{-g} R \phi^2$$

to the Lagrangian. For which value  $\xi$  does the trace vanish?

- d) (2 points) Compute the field equations for the scalar  $\phi$  for the matter Lagrangian used in the last task and show that  $S_{R\phi^2}$  induces a mass correction for the  $\phi$  proportional to the curvature R.
- e) (2 points) Show that for the value of  $\xi$  computed in task c, the action for  $\phi$  is invariant under local *conformal transformations*

$$g_{\mu\nu} \to g'_{\mu\nu} = \Omega^{-2}(x)g_{\mu\nu}$$

## Exercise 4.2: Coordinates of $AdS_{d+1}$

We approach Anti-de Sitter space through the hypersurface

$$-L^{2} = \tilde{\eta}_{MN} X^{M} X^{N} = -(X^{d+1})^{2} - (X^{0})^{2} + \sum_{i+1}^{d} (X^{i})^{2} , \qquad (3)$$

(26 points)

where  $X \in \mathbb{R}^{d,2}$  with  $ds^2 = \tilde{\eta}_{MN} X^M X^N$ . In the following we will use two different parameterizations. The first one, is called *global coordinates* with  $(\rho, \tau, \Omega_i)$ ,

$$X^{d+1} = L \cosh \rho \sin \tau$$
$$X^0 = L \cosh \rho \cos \tau$$
$$X^i = L \sinh \rho \Omega_i.$$

We have encountered them in the lecture and remember that i = 1, ..., d and  $\sum_{i=1}^{d} \Omega_i^2 = 1$ .

- a) (2 points) Compute the induced metric  $ds^2 = g_{\mu\nu}dx^{\mu}dx^{\nu}$  in global coordinates. *Hint:* Revisit the computation from the two-sphere, I presented in the  $2^{nd}$  lecture.
- b) (2 points) Replace  $\rho$  by  $r = L \sinh \rho$  and show that the metric can be written in the form

$$ds^{2} = -H(r)dt^{2} + H(r)^{-1}dr^{2} + r^{2}d\Omega_{d-1}^{2}$$

where  $d\Omega_{d-1}^2$  is the metric of the unit (d-1)-sphere,  $S^{d-1}$ .

c) (3 points) The Poincaré patch coordinates  $(x^{\mu}, u), \mu = 0, \dots, d-1$ , are defined by

$$X^{d+1} + X^d = u$$
$$-X^{d+1} + X^d = v$$
$$X^{\mu} = \frac{u}{L} x^{\mu}.$$

Using (3), eliminate v in terms of u and  $x^{\mu}$  and show that the induced metric for  $(u, x^{\mu})$  takes the compact form

$$ds^2 = L^2 \frac{du^2}{u^2} + \frac{u^2}{L^2} dx^{\mu} dx_{\mu}.$$

Finally, introduce  $z = \frac{L^2}{u}$  and show that the metric takes the very simple form

$$ds^{2} = \frac{L^{2}}{z^{2}} \left( dz^{2} + dx^{\mu} dx_{\mu} \right) .$$
 (4)

d) (2 points) Redraw the embedding of  $AdS_2$ , in  $\mathbb{R}^{1,2}$  presented in the lecture. Which part of this AdS spacetime is not covered by Poincaré coordinates? *Hint: z only takes positive values. Why?* 

## Exercise 4.3: Curvature of AdS and the cosmological constant

In the last problem we have found a very simple form from the metric of Anti-de Sitter spacetime in Poicaré coordinates given in (4).

- a) (3 points) Compute the Ricci tensor  $R_{\mu\nu}$  and the Ricci scalar R for this metric.
- b) (2 points) Show with the results from the last task that the AdS spacetime solves the vacuum Einstein field equations (1) (with vanishing energy-momentum). Determine the required cosmological constant.