Gauge/Gravity Duality, Winter 2024/25

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2. Abelian and non-Abelian gauge theory (18 points)

To be discussed on Friday, 11^{th} October, 2024 in the tutorial. Please indicate your preferences until Sunday, 06/10/2024, 21:00:00 on the website.

Exercise 2.1: Global symmetries

Consider the action for a U(1) gauge theory

$$S = -\frac{1}{4g^2} \int d^4 x F_{\mu\nu} F^{\mu\nu} \quad \text{with} \quad F_{\mu\nu} = 2\partial_{[\mu} A_{\nu]} \,. \tag{1}$$

Under a Poincaré transformation $x^{\mu} \to x'^{\mu} = \Lambda^{\mu}{}_{\nu}x^{\nu} + a^{\mu}$, the gauge field A_{μ} transforms as

$$A^{\mu}(x) \to A'^{\mu}(x) = \Lambda^{\mu}{}_{\nu}A^{\nu}(\Lambda^{-1}x - a).$$
 (2)

a) (2 points) Show that the action (1) is invariant under (2). Moreover, derive the infinitesimal transformation law

$$\delta A_{\mu} = \omega_{\nu\lambda} x^{\nu} \partial^{\lambda} A_{\mu} - \omega_{\mu}{}^{\nu} A_{\nu} - a^{\nu} \partial_{\nu} A_{\mu}$$

by using $\Lambda^{\mu}{}_{\nu} = \delta^{\mu}_{\nu} + \omega^{\mu}{}_{\nu}$ with $\omega_{\mu\nu} = -\omega_{\nu\mu}$.

b) (3 points) Derive the conserved current, $\Theta_{\mu\nu}$ satisfying $\partial^{\mu}\Theta_{\mu\nu} = 0$, corresponding to infinitesimal translations with the parameter a^{μ} . You will note that it is not symmetric. However, we can construct a symmetric tensors which is still conserved, namely

$$T_{\mu\nu} = \Theta_{\mu\nu} + \partial^{\lambda} f_{\lambda\mu\nu} \, ,$$

by choosing $f_{\mu\nu\lambda} = -f_{\nu\mu\lambda}$ appropriately. Show that $f_{\lambda\mu\nu} = -\frac{1}{g^2}A_{\nu}F_{\mu\lambda}$ works and will give the canonical energy-momentum tensor

$$T_{\mu\nu} = \frac{1}{g^2} \left(F_{\mu\lambda} F_{\nu}{}^{\lambda} - \frac{1}{4} \eta_{\mu\nu} F_{\alpha\beta} F^{\alpha\beta} \right)$$

of electrodynamics.

c) (2 points) Derive the conserved current for Lorentz transformations with the parameter $\omega_{\mu\nu}$. It is of the form

$$N^{\nu\kappa\lambda} = x^{\kappa}T^{\lambda\nu} - x^{\lambda}T^{\kappa\nu}$$

Check the this current is conserved and therefore satisfies $\partial^{\mu} N_{\mu\nu\lambda} = 0$ for any solutions of the field equations.

Exercise 2.2: Yang-Mills theory

Now, we look at the action of a non-Abelian gauge theory

$$S = -\frac{1}{2g^2} \int \mathrm{d}^4 x \, \mathrm{Tr}(F_{\mu\nu}F^{\mu\nu}) \tag{3}$$

with the field strength $F_{\mu\nu} = F^a_{\mu\nu}T_a = 2\partial_{[\mu}A_{\nu]} + i[A_{\mu}, A_{\nu}]$ and the gauge potential $A_{\mu} = A^a_{\mu}T_a$.

a) (2 points) Assuming real structure coefficient $f_{ab}{}^c$ in the commutator

$$[T_a, T_b] = i f_{ab}{}^c T_c \,,$$

show that the generators of the underlying Lie algebra T_a can be chosen to be hermitian. Is it possible to generalize this choice? If yes, how? Argue why the pairing

$$Tr(T_a T_b) = \kappa_{ab}$$

is real, too.

b) (2 points) Show that the field strength transforms as

$$F_{\mu\nu} \to F'_{\mu\nu} = U F_{\mu\nu} U^{\dagger}$$

assuming the transformation

$$A_{\mu} \to A'_{\mu} = U A_{\mu} U^{\dagger} - i U \partial_{\mu} U^{\dagger}$$

with $U = \exp[i\alpha^a(x)T_a]$ and real, coordinate dependent parameters $\alpha^a(x)$.

c) (2 points) Show that the equations of motion for the action (3) can be written as

$$D_{\mu}F^{\mu\nu} = 0$$
 with $D_{\mu}F^{\nu\rho} = \partial_{\mu}F^{\nu\rho} + i[A_{\mu}, F^{\nu\rho}].$

d) (2 points) Verify that the Bianchi identity

$$3D_{[\mu}F_{\nu\rho]} = 0$$

automatically holds.

e) (3 points) As we consider here four dimensions, a term of the form

$$S^{\text{top}} = \int d^4x \frac{\vartheta}{16\pi^2} \operatorname{Tr}(F_{\mu\nu}\tilde{F}^{\mu\nu})$$
(4)

may also be added to the action, with $\tilde{F}^{\mu\nu} = \frac{1}{2} \epsilon^{\mu\nu\rho\sigma} F_{\rho\sigma}$ denoting the dual field strength tensor (dual under Hodge duality). Here, we encounter the totally antisymmetric tensors $\epsilon^{\mu\nu\rho\sigma}$ in four dimensions. It is normalized such that $\epsilon^{0123} = -1$. Show that (4) is also gauge invariant but does not contribute to the equations of motion.