

## 7. Iteration, bifurcation and chaos

Chaos already in 5<sup>th</sup> lecture. But complicated setup, i.e. double pendulum, ODE, etc.

Is there an even simpler setup with similar behavior?

### 7.1. Logistic map

$$f(x) = r \cdot x(1-x)$$

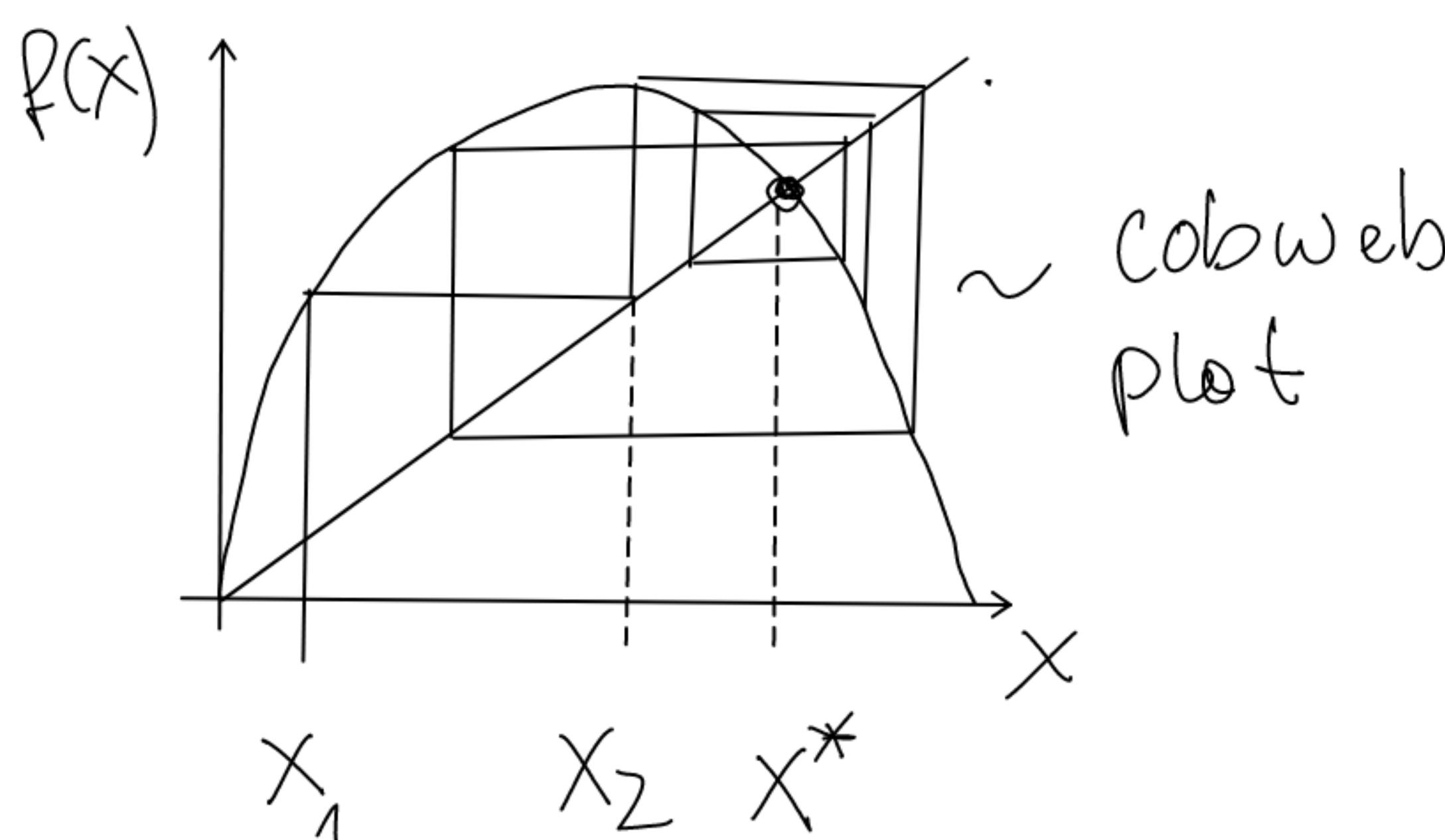
rabbit population:

$x < \frac{1}{2}$  reproduction

$x > \frac{1}{2}$  starvation

$r$  ≈ rate of reproduction

Iteration  $x_{i+1} = f(x_i)$  can be visualized by



and has fixed point.

When does it converge?

Remember Banach's theorem.

$0 < r < 1$  :  $f(x)$  is contraction on  $A = [0,1]$

(boring)  $\Rightarrow$  All  $x_0 \in A$  converge to  $x^* = 0$

for  $r > 1$  :  $x = rx(1-x) \rightarrow x \underbrace{(-rx + r - 1)}_{=0} = 0$

$$x_1^* = 0 \quad \& \quad x_2^* = 1 - \frac{1}{r}$$

If we are close to a fixed point:

$$x_{i+1} = f(x^* + \varepsilon_i) \approx f(x^*) + \underbrace{f'(x^*)}_{\varepsilon_i}_{\varepsilon_{i+1}}$$

$\Rightarrow \varepsilon_{i+1} = f'(x^*) \varepsilon_i \rightarrow$

$ f'(x^*)  < 1 \rightarrow$ stable (contraction)
$ f'(x^*)  > 1 \rightarrow$ unstable

for the logistic map  $f'(x_1^*) = r$

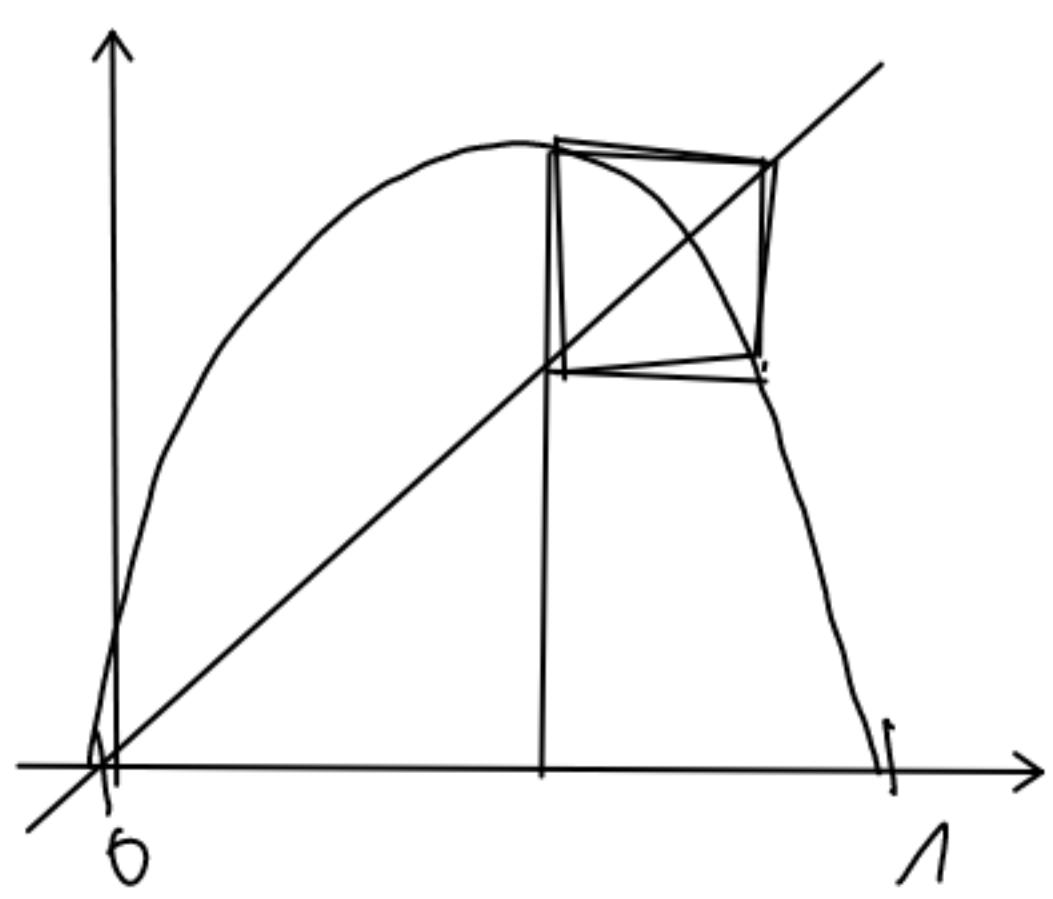
$$f'(x_2^*) = 2 - r$$

$\Rightarrow 1 < r < 3 \quad x_1^* = 0 \quad$  unstable and  
 $x_2^* = 1 - 1/r \quad$  stable

Observation: iteration converges to stable fixed point.

### 3.2. Bifurcation and periodic orbits

For  $r=3+\varepsilon$ , we see periodic orbits like



with period 2.

$\Leftrightarrow$  fixed point of  $f^2(x) = g(x)$

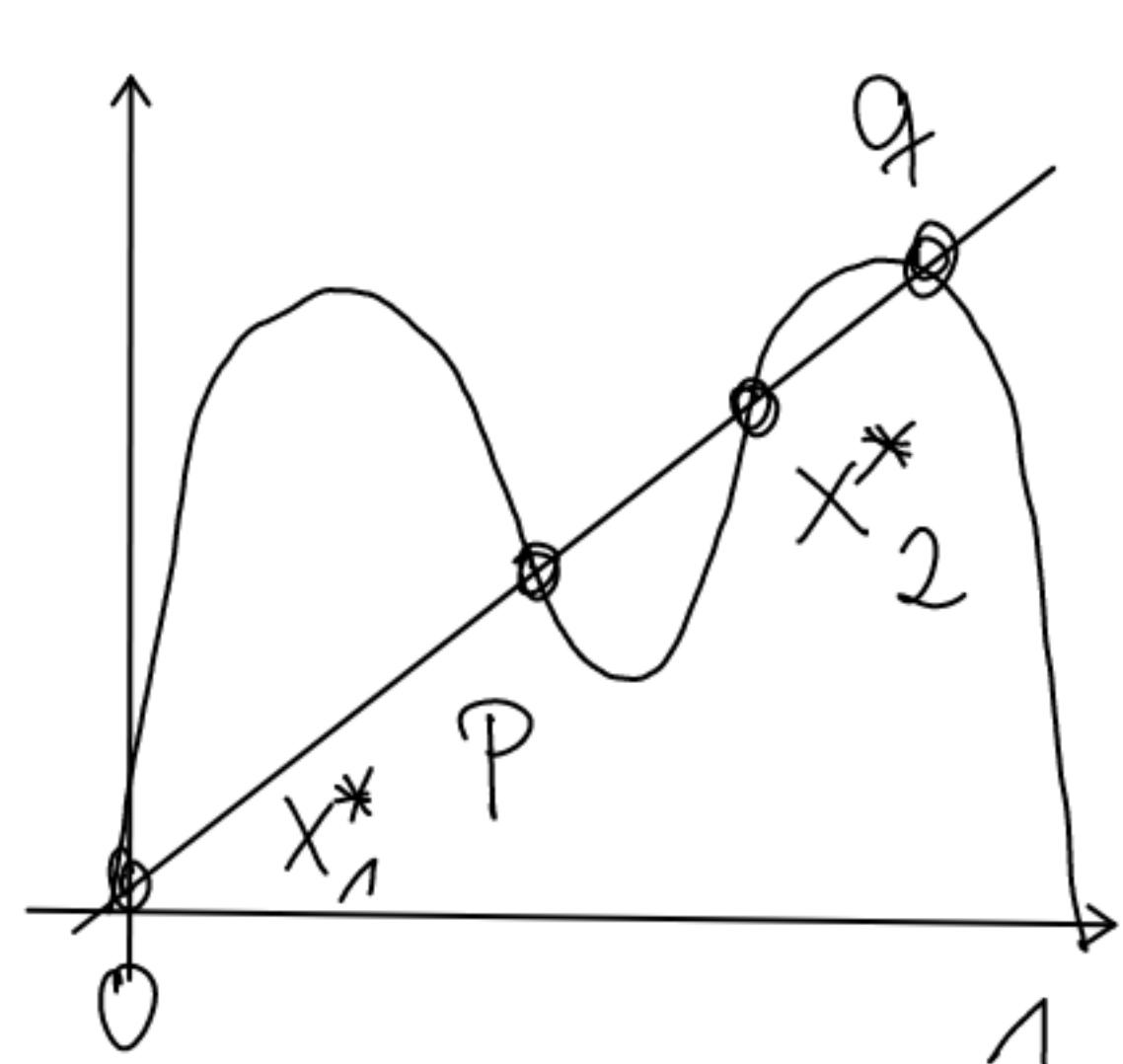
two new fixed points, p & q

$$x_1^* = 0, \quad x_2^* = 1 - 1/r \quad (\text{unstable for } r > 3)$$

$$p, q = \frac{r+1 \pm \sqrt{(r-3)(r+1)}}{2r}$$

stable  $3 < r < 3.449\dots$

When they become unstable, we go to



$h(x) = g(x)^2 = f(x)^4 \rightarrow$  periodic orbit with period 4 ...

$r <$	Period	
$r_1 = 3$	2	
$r_2 = 3,449..$	4	
$r_3 = 3,54409..$	8	
$\vdots$	$\vdots$	
$r_\infty = 3,569946..$	$\infty$	

bifurcation

Feigenbaum - constant

new unstable fixed point for  $f^{2^n}(x)$  and two new stable fixed points around it.

### 3.3. Chaos

$r > 3,5699..$  no periodic orbit  $\rightarrow$  chaos

still some structure, i.e. window with period 3

for  $3,8284 < r < 3,8415 \rightarrow$  check for

$$r = 3,8282 \quad \& \quad x_0 = \frac{1}{2}$$

Question: How to see this structure?

~~Do a lot of iterations ( $\sim 1000$ ) starting close to  $x^* = 0$ , i.e.  $x_0 = \epsilon = 10^{-5}$  and plot the last 100 on the y-axis. Repeat for different  $r$  on the x-axis.~~

$\rightsquigarrow$  bifurcation diagram

### 3.4. Self-similarity and universality

unimodal

"zooming" into the diagram  
we see the same structures

all smooth concave maps  
with one maximum on  
 $[0, 1]$  have same  $\delta$

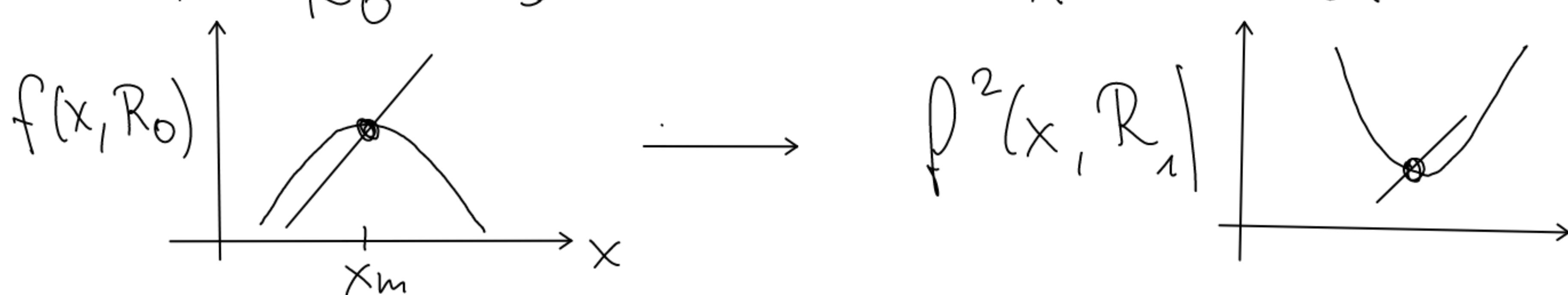
Why? RG-flow

① Superstable fixed point for  $f'(x^*) = 0$

remember from last lecture at least quadratic convergence (Newton's method with  $p=2$ )

② Take  $f(x)$  with maximum at  $x_m$  and choose

$r = R_0$  such that  $x^* = x_m$ .



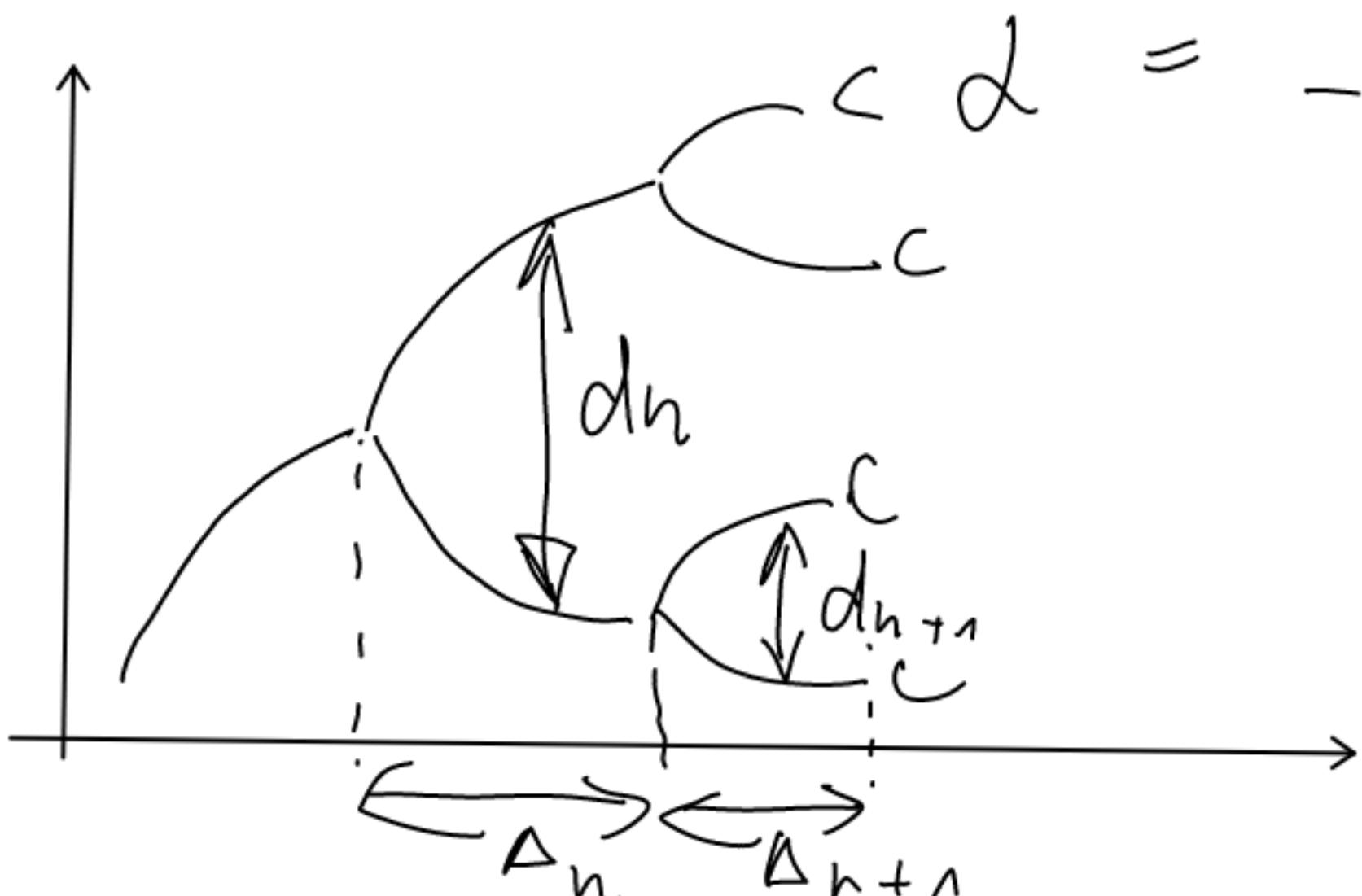
③ Rescale the coordinates to get

$$f(x, R_0) \approx \lambda f^2\left(\frac{x}{\lambda}, R_1\right)$$

④ Alternatively  $r = R_i$  such that  $x_m$  is an orbit of period 2

converges to  $g_\infty(x) = \lim_{n \rightarrow \infty} \lambda^n f^{2^n}\left(\frac{x}{\lambda^n}, R_{n+i}\right)$

for  $i \rightarrow \infty$   $g_\infty(x) = \lambda g_\infty\left(\frac{x}{\lambda}\right)$  fixed point function



$$\delta = \lim_{n \rightarrow \infty} \frac{\Delta_n}{\Delta_{n+1}}$$

$$\lambda = \lim_{n \rightarrow \infty} \frac{d_n}{d_{n+1}}$$