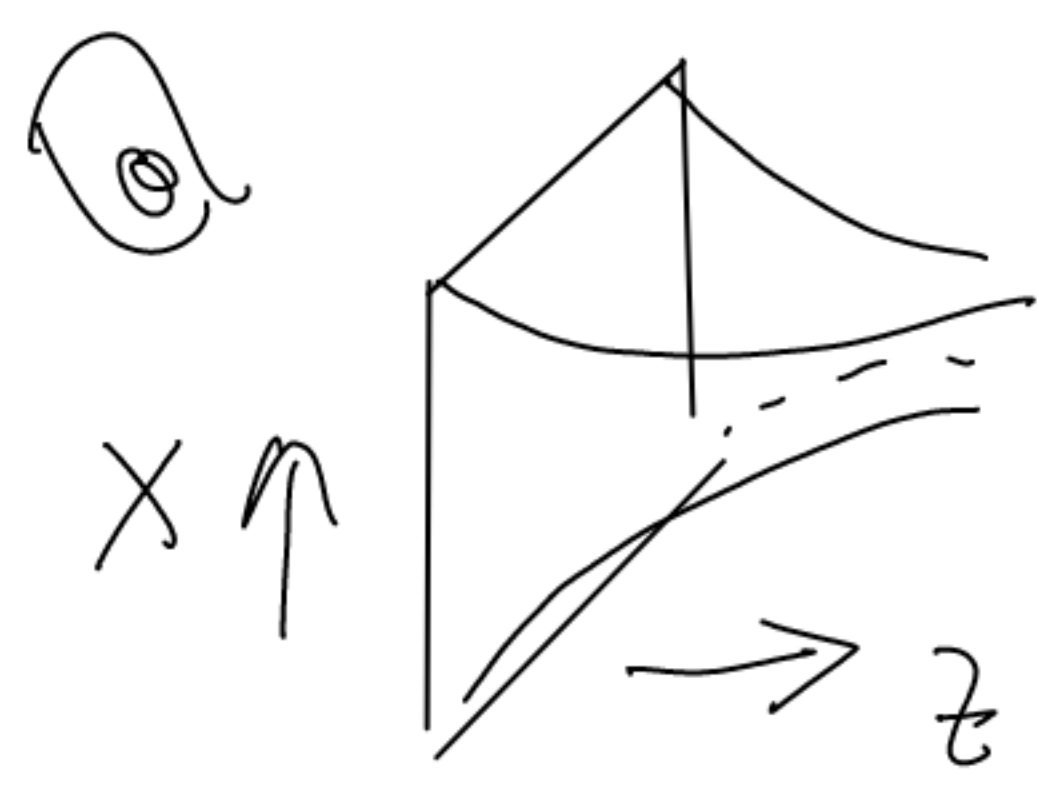


11. Holographic renormalization



z in AdS seems to be related to an energy (cut-off) scale.
 \rightarrow renormalization (\nearrow 4.4.)

scalars from last lecture in Fefferman-Graham coord.

with $ds^2 = L^2 \left(\frac{ds^2}{4s^2} + \frac{1}{s} \delta_{\mu\nu} dx^\mu dx^\nu \right)$ boundary @ $s=0$

ansatz for eom solution: $\phi(x, s) = s^{(d-\Delta)/2} \bar{\phi}(x, s)$

with $\bar{\phi}(x, s) = \phi_{(0)}(x) + s \phi_{(2)}(x) + s^2 \phi_{(4)}(x) + \dots$

$$\Rightarrow \phi_{(2n)} = \frac{1}{2n(2\Delta - d - 2n)} \square_0 \phi_{(2n-2)} \quad \delta^{\mu\nu} \partial_\mu \partial_\nu$$

⚡ Problem for $2\Delta - d - 2k = 0$ for any $k \in \mathbb{N}$

i.e. $k=1$ with $\Delta = \frac{1}{2}(d-2) \Rightarrow$ modify ansatz

- $\bar{\phi}(x, s) = \phi_{(0)} + s (\phi_{(2)} + \ln s \chi_{(2)})$ with

$$\chi_{(2)} = -\frac{1}{4} \square_0 \phi_{(0)},$$

$\phi_{(2)}$ more complicated, non-locally (∞ ∂ 's) related to $\phi_{(0)}$

for general $k \in \mathbb{N}$ $\chi_{(2k)} = -\frac{1}{2^{2k} \Gamma(k) \Gamma(k+1)} (\square_0)^k \phi_{(0)}$

like before into the action, only boundary contributes

$$S_{\text{reg}} = -\frac{C}{2} \int d^d x \sqrt{g} g^{\beta\beta} \phi \partial_s \phi \Big|_{s=\epsilon}$$

$$= CL^{d-1} \int d^d x \left(\varepsilon^{-\Delta + \frac{d}{2}} a_{(0)} + \varepsilon^{-\Delta + \frac{d}{2} + 1} a_{(1)} + \dots - \ln \varepsilon a_{(2\Delta-d)} \right)$$

$$= -\frac{1}{2} (d-\Delta) \phi_{(0)}^2 = -(d-\Delta+1) \phi_{(0)} \phi_{(2)} \dots$$

for $\Delta > d/2$ S_{reg} diverges for $\varepsilon \rightarrow 0$

~~$\frac{\varepsilon}{c}$~~ compensate by adding counter terms to the action.

$$\gamma_{\mu\nu} = L^2 \delta_{\mu\nu} / \varepsilon$$

$$S_{\text{ct}} = \frac{C}{L} \int d^d x \sqrt{\gamma} \left(\frac{d-\Delta}{2} \phi^2(x, \varepsilon) + \frac{1}{2(2\Delta-d-2)} \phi(x, \varepsilon) \square \phi(x, \varepsilon) \right)$$

here for $d/2 < \Delta < d/2 + 1$, for larger Δ more ct

now the renormalized action $S_{\text{ren}} = \lim_{\varepsilon \rightarrow 0} S_{\varepsilon}$
with $S_{\varepsilon} = S_{\text{reg}} + S_{\text{ct}}$

plugging in the solution for $\phi(x)$ we get the finite contributions

$$S_{\text{reg}} = -\frac{dC L^{d-1}}{2} \int d^d x \phi_{(0)}(x) \phi_{(2)}(x) + \dots$$

$$S_{\text{ct}} = CL^{d-1} (d-\Delta) \int d^d x \phi_{(0)}(x) \phi_{(2)}(x) + \dots$$

both together give

$$S_{\varepsilon} = \frac{CL^{d-1}}{2} (d-2\Delta) \frac{\Gamma(\Delta)}{\pi^{d/2} \Gamma(\Delta - \frac{d}{2})} \int d^d x \int d^d y \frac{\phi_{(0)}(x) \phi_{(0)}(y)}{(x-y)^{2\Delta}}$$

and we again get

$$\langle \theta(x) \theta(y) \rangle = CL^{d-1} (2\Delta-d) \frac{\Gamma(\Delta)}{\pi^{d/2} \Gamma(\Delta - \frac{d}{2})} \frac{1}{(x-y)^{2\Delta}}$$

Remark: Same method of constructing counter

terms also applies to the metric.