


with this ansatz we find \rightarrow see EX 11.2

$$R(G) = R(g) - \frac{1}{4} \kappa^2 F^{ij} F_{ij} - \frac{2}{\kappa} \nabla^i \nabla_i \kappa$$

$$F_{ij} = 2 \partial_{[i} A_{j]} \leftarrow \begin{array}{l} \text{electro-magnetic} \\ \text{field strength} \end{array}$$

$$\hookrightarrow S = \int d^{D-1} x \sqrt{-g} \left(R - (\partial \ln \kappa)^2 - \frac{\kappa^2}{4} F^{ij} F_{ij} \right)$$

GR + ED unified

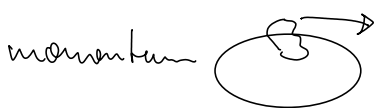
\hookrightarrow massless scalar \rightarrow not observed 

Same in more sophisticated compactifications
 $\hat{=}$ moduli stabilisation problem

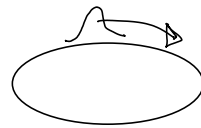
\leadsto use H-flux $\hat{=}$ flux compactifications
 \leadsto research

12. Closed string on torus & T-duality

Motivation: strings \neq point particles



=



BUT

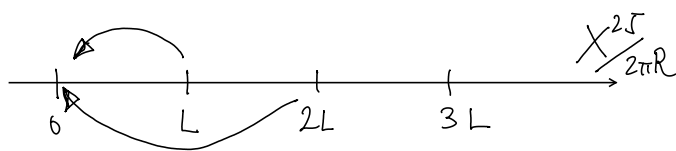
\times

Important for i.e. Kaluza-Klein. Implications?

12.1. Closed string on a circle

remember: $D=26 \leadsto$ compactify X^{25}

$$X^{25} \sim X^{25} + 2\pi R \cdot L \quad L \in \mathbb{Z}$$



$$S^1 = \mathbb{R} / (2\pi R \mathbb{Z}) \quad \text{and therefore}$$

$$\boxed{X^{25}(\sigma + 2\pi, \tau) = X^{25}(\sigma) + 2\pi R L}$$

→ into mode expansion

$$X^{25}(\sigma, \tau) = \underbrace{X^{25} + d' p^{25} \cdot \tau + L R_0}_{\text{zero modes}} + i \sqrt{\frac{\alpha'}{2}} \sum_{n \neq 0} \text{oscillators}$$

Zero modes with $[X^{25}, p^{25}] = i$
Wave function on S^1 single-valued

$$\longrightarrow \boxed{P = \frac{M}{R} \Delta} \quad M \in \mathbb{Z} \quad \text{momentum quantum number}$$

$$= X_L^{25}(\tau + \sigma) + X_R^{25}(\tau - \sigma) \quad (\text{know from KK})$$

↳ Left- and Right moving modes

$$X_L^{25}(\tau + \sigma) = \frac{1}{2} (X^{25} + c) + \underbrace{\frac{\alpha'}{2} \left(\frac{M}{R} + \frac{LR}{\alpha'} \right)}_{P_L^{25}} (\tau + \sigma) + \text{OSC.}$$

$$X_R^{25}(\tau - \sigma) = \frac{1}{2} (X^{25} - c) + \underbrace{\frac{\alpha'}{2} \left(\frac{M}{R} - \frac{LR}{\alpha'} \right)}_{P_R^{25}} (\tau - \sigma) + \text{OSC.}$$

with mass

$$\alpha' m_L^2 = 2(L_0 - 1) = \frac{\alpha'}{2} \left(\frac{M}{R} + \frac{LR}{\alpha'} \right)^2 + 2(N_L - 1)$$

$$\alpha' m_R^2 = 2(L_0 - 1) = \frac{\alpha'}{2} \left(\frac{M}{R} - \frac{LR}{\alpha'} \right)^2 + 2(N_R - 1)$$

rest frame in $\mathbb{R}^{1,24}$ uncompactified space

remember Level matching

$$L_0 |\psi\rangle = \bar{L}_0 |\psi\rangle$$

physical state

$$\Rightarrow m_L^2 = m_R^2 \quad \text{or}$$

$$\boxed{N_R - N_L = M \cdot L}$$

with mass

$$\alpha' m^2 = \alpha' (m_L^2 + m_R^2) = \alpha' \frac{M^2}{R^2} + \frac{1}{\alpha'} L^2 R^2 + 2(N_L + N_R - 2)$$

momentum
energy of and
winding modes

Spectrum:

- ① $|0\rangle$ $\alpha' m^2 = -4 \Rightarrow$ tachyon
- ② $|G^{\mu\nu}\rangle = \alpha_{-1}^\mu \bar{\alpha}_{-1}^\nu |0\rangle \Rightarrow$ massless spin 2
metric, B-field & dilaton in $\mathbb{R}^{1,24}$
- ③ $|V_1^M\rangle = \alpha_{-1}^M \bar{\alpha}_{-1}^{25} |0\rangle \Rightarrow$ massless spin 1
 $|V_2^M\rangle = \alpha_{-1}^{25} \bar{\alpha}_{-1}^M |0\rangle$ two vectors
- ④ $|\phi\rangle = \alpha_{-1}^{25} \bar{\alpha}_{-1}^{25} |0\rangle \Rightarrow$ massless scalar

for $M=L=0$ we recover fields from KK

- BUT much more excited states (in KK only $L=0$)
- \mathbb{Z}_2 symmetry of spectrum

$$R \Leftrightarrow \frac{\alpha'}{R}, \quad L \Leftrightarrow M \quad \text{called } T\text{-duality}$$

$T =$ Target space or Torus

- not only spectrum but also modes with

$$(P_L^{25}, P_R^{25}) \rightarrow (P_L^{25}, -P_R^{25}) \quad \text{and osc. such that}$$

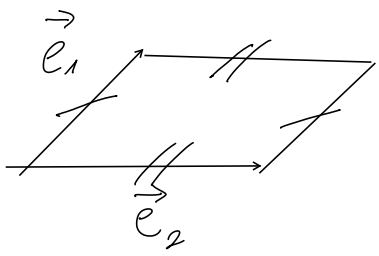
$$(X_L^{25}, X_R^{25}) \rightarrow (X_L^{25}, -X_R^{25})$$

full sym. of closed string!

12.2 Closed string on D-torus

$$X^I \sim X^I + 2\pi \sum_{i=1}^D n^i e_i^I = L^I$$

torus directions



$T^D = \mathbb{R}^D / 2\pi \Lambda_D$
 lattice spanned by \vec{e}_i
 Torus

again $[X^I, P_J] = i \delta^I_J$, $P_I = m_i \vec{e}_I^*$ quantised
 \vec{e}_I^* = dual lattice with

$$e_i^I e_I^j = \delta_i^j \quad \text{and} \quad e_i^I e_j^* = \delta_{ij}^I$$

⋮

mass formula $\alpha' m^2 = 2(N_L + N_R - 2) + \sum_{i,j=1}^D (\alpha' m_i g^{ij} m_j + \frac{1}{\alpha'} n^i g_{ij} n^j)$

$g_{ij} = e_i^I e_j^J \delta_{IJ}$

$g^{ij} = e_I^* e_j^* \delta^{IJ}$

and

level matching $N_R - N_L = P_I L^I = \sum_{i=1}^D m_i n^i$

$$= \frac{1}{2} N^{\hat{I}} \eta_{\hat{I}\hat{J}} N^{\hat{J}}$$

with $N^{\hat{I}} = \begin{pmatrix} n^i \\ m_i \end{pmatrix}$ and $\eta_{\hat{I}\hat{J}} = \begin{pmatrix} 0 & \delta_i^j \\ \delta_i^j & 0 \end{pmatrix}$

invariant metric of the Lie group $O(D,D)$

similar $\alpha' m^2 = \text{osc} + N^{\hat{I}} \mathcal{H}_{\hat{I}\hat{J}} N^{\hat{J}}$

with $\mathcal{H}_{\hat{I}\hat{J}} = \begin{pmatrix} \frac{1}{\alpha'} g_{ij} & 0 \\ 0 & \alpha' g^{ij} \end{pmatrix} \hat{=} \text{generalised metric}$

T-duality $\hat{=} O(D,D, \mathbb{Z})$ action

non-linearly realised on g_{ij} and $B_{ij} \Rightarrow$ Buscher rules