

# The Generalized Geometry of $\alpha'$ -corrections

Falk Hassler

Based on work in progress with

Daniel Butter, Achilles Gitsis & Eric Lescano



Uniwersytet  
Wrocławski



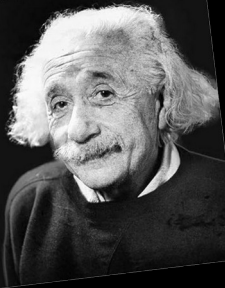
# The Problem

- Einstein-Hilbert action is not renormalizable in  $d > 2$   $\longrightarrow$  only EFT

$$S = \int dx^d \sqrt{-g} (R + a_1 R^2 + a_2 R_{ij} R^{ij} + \dots)$$

Do not worry about  
your difficulties in  
mathematics. I can  
assure you mine are  
still greater.

ALBERT EINSTEIN



# The Problem

- Einstein-Hilbert action is not renormalizable in  $d > 2$   $\longrightarrow$  only EFT

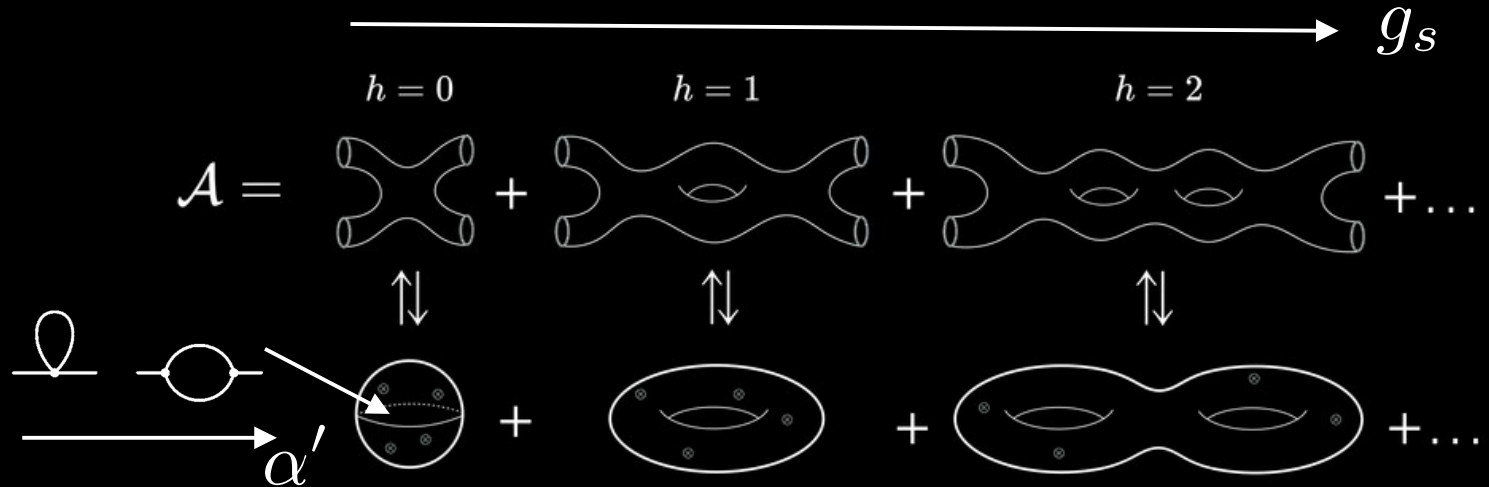
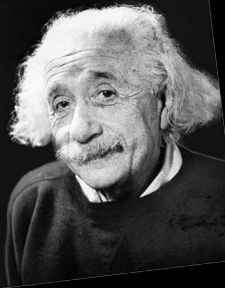
$$S = \int dx^d \sqrt{-g} (R + a_1 R^2 + a_2 R_{ij} R^{ij} + \dots)$$

Question: How do we obtain all the coefficients ?

String Theory

Do not worry about your difficulties in mathematics. I can assure you mine are still greater.

ALBERT EINSTEIN



# NS/NS-sector @ leading order in $\alpha'$

$$S = \int dx^d \sqrt{-g} e^{-2\phi} \left( R + 4(\partial\phi)^2 - \frac{1}{12} \tilde{H}^2 + \frac{a}{8} R_{ija}^{(-)b} R^{(-)ij}{}_b{}^a + \frac{b}{8} R_{ija}^{(+b)} R^{(+ij)}{}_b{}^a + \dots \right)$$

$\tilde{H}_{ijk} = H_{ijk} - \frac{3}{2}a\Omega_{ijk}^{(-)} + \frac{3}{2}b\Omega_{ijk}^{(+)}$

$a = -\alpha, b = 0$	heterotic
$a = b = -\alpha'$	bosonic
$a = b = 0$	type II

# NS/NS-sector @ leading order in $\alpha'$

$$S = \int dx^d \sqrt{-g} e^{-2\phi} \left( R + 4(\partial\phi)^2 - \frac{1}{12} \tilde{H}^2 + \frac{a}{8} R_{ija}^{(-)b} R^{(-)ij}{}_b{}^a + \frac{b}{8} R_{ija}^{(+b)} R^{(+ij)}{}_b{}^a + \dots \right)$$

$\tilde{H}_{ijk} = H_{ijk} - \frac{3}{2}a\Omega_{ijk}^{(-)} + \frac{3}{2}b\Omega_{ijk}^{(+)}$

$a = -\alpha, b = 0$	heterotic
$a = b = -\alpha'$	bosonic
$a = b = 0$	type II

- 3 coefficients for terms with 2
  - 8 coefficients for terms with 4
  - 60 coefficients for terms with 6
- } derivatives

# NS/NS-sector @ leading order in $\alpha'$

$$S = \int dx^d \sqrt{-g} e^{-2\phi} \left( R + 4(\partial\phi)^2 - \frac{1}{12} \tilde{H}^2 \right. \\ \left. + \frac{a}{8} R_{ija}^{(-)b} R^{(-)ij}{}_b{}^a + \frac{b}{8} R_{ija}^{(+b)} R^{(+ij)}{}_b{}^a + \dots \right)$$

$\tilde{H}_{ijk} = H_{ijk} - \frac{3}{2}a\Omega_{ijk}^{(-)} + \frac{3}{2}b\Omega_{ijk}^{(+)}$

$a = -\alpha, b = 0$	heterotic
$a = b = -\alpha'$	bosonic
$a = b = 0$	type II

- 3 coefficients for terms with 2
  - 8 coefficients for terms with 4
  - 60 coefficients for terms with 6
- } derivatives



Too many terms. Nothing is known about  $> 8$  derivatives.

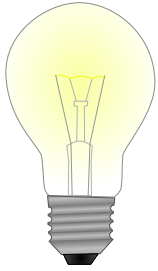


# A better approach:

Leverage symmetry to decrease number of possible terms.

Like diffeomorphisms, gauge-transformations and:

- SUSY
- Extended Generalized Lorentz Symmetry (today)
- ...



# A better approach:

Leverage symmetry to decrease number of possible terms.

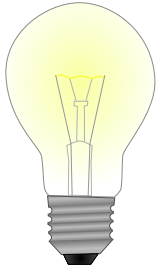
Like diffeomorphisms, gauge-transformations and:

- SUSY
- Extended Generalized Lorentz Symmetry (today)
- ...

generalized frame

$$E_A^I = \begin{pmatrix} e_a^i & e_a^j B_{ji} \\ 0 & e^a_i \end{pmatrix}$$



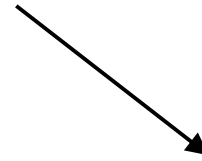


# A better approach:

Leverage symmetry to decrease number of possible terms.

Like diffeomorphisms, gauge-transformations and:

- SUSY
- Extended Generalized Lorentz Symmetry (today)
- ...



generalized frame

invariant under  $O(d) \times O(d) \subset O(d,d)$

$$E_A^I = \begin{pmatrix} e_a^i & e_a^j B_{ji} \\ 0 & e^a_i \end{pmatrix}$$

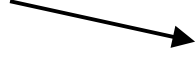
$$\eta_{AB} = \begin{pmatrix} 0 & \delta_{\alpha}^{\beta} \\ \delta_{\beta}^{\alpha} & 0 \end{pmatrix} \quad \mathbb{H}_{AB} = \begin{pmatrix} \delta_{ab} & 0 \\ 0 & \delta^{ab} \end{pmatrix}$$

# Leading Symmetries and Action

$$\delta E^A_M = \mathbb{L}_\xi E^A_M + \Lambda^A_B E^B_M, \quad \Lambda^A_B \in O(d) \times O(d)$$



generalized Lie derivative



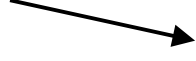
generalized Lorentz transformation

- 1) diffeomorphisms (gravity)
- 2) gauge transformation

transformation of fermions

# Leading Symmetries and Action

$$\delta E^A_M = \mathbb{L}_\xi E^A_M + \Lambda^A_B E^B_M, \quad \Lambda^A_B \in \text{O}(d) \times \text{O}(d)$$



generalized Lie derivative

generalized Lorentz transformation

1) diffeomorphisms (gravity)

2) gauge transformation

transformation of fermions

generalized flux  $F_{ABC} = 3D_{[A} E_B^I E_{C]I}$  with  $D_A = E_A^i \partial_i$

$$F_A = D_A d - \partial_i E_A^i \quad d = -\frac{1}{2} \log(-g) + \phi$$

# Leading Symmetries and Action

$$\delta E^A_M = \mathbb{L}_\xi E^A_M + \Lambda^A_B E^B_M, \quad \Lambda^A_B \in O(d) \times O(d)$$



generalized Lie derivative



generalized Lorentz transformation

1) diffeomorphisms (gravity)

2) gauge transformation

transformation of fermions

generalized flux  $F_{ABC} = 3D_{[A} E_B^I E_{C]I}$  with  $D_A = E_A^i \partial_i$

$$F_A = D_A d - \partial_i E_A^i \quad d = -\frac{1}{2} \log(-g) + \phi$$

$$S = \int dx^d e^{-2d} \mathcal{R}$$

one unique invariant

$$\mathcal{R}(F_{ABC}, F_A, D_A, H_{AB})$$

# We need a Factory for Invariants...

## Symmetries

- gen. diff
- gen. Lorentz



## Invariants

$\mathcal{R}, \dots$

Poláček-Siegel construction \*

\*) for mathematicians: symplectic reduction of Courant algebroids

# We need a Factory for Invariants...

## Symmetries

- gen. diff
- gen. Lorentz



## Invariants

$\mathcal{R}, \dots$

Poláček-Siegel construction \*

## Covariant derivative:

$$\nabla_A E_B^M = E_A^N \partial_N E_B^M + \Omega_{AB}^C E_C^M - E_A^N \Gamma_{NL}^M E_B^L$$

gen. spin and affine connection, related by

$$\nabla_A E_B^M = 0$$

\*) for mathematicians: symplectic reduction of Courant algebroids

# We need a Factory for Invariants...

## Symmetries

- gen. diff
- gen. Lorentz



## Invariants

$\mathcal{R}, \dots$

Poláček-Siegel construction \*

## Covariant derivative:

$$\nabla_A E_B^M = E_A^N \partial_N E_B^M + \Omega_{AB}^C E_C^M - E_A^N \Gamma_{NL}^M E_B^L$$

gen. spin and affine connection, related by

$$\nabla_A E_B^M = 0$$

## Curvature and torsion???:

$$[\nabla_A, \nabla_B] V^C = R_{ABD}^C V^D + T_{AB}^D \nabla_D V^C$$

\*) for mathematicians: symplectic reduction of Courant algebroids

# We need a Factory for Invariants...

## Symmetries

- gen. diff
- gen. Lorentz



## Invariants

$\mathcal{R}, \dots$

Poláček-Siegel construction \*

## Covariant derivative:

$$\nabla_A E_B^M = E_A^N \partial_N E_B^M + \Omega_{AB}^C E_C^M - E_A^N \Gamma_{NL}^M E_B^L$$

gen. spin and affine connection, related by

$$\nabla_A E_B^M = 0$$

## Curvature and torsion???:

$$[\nabla_A, \nabla_B] V^C = R_{ABD}^C V^D + T_{AB}^D \nabla_D V^C$$



$R_{ABC}^D$  &  $T_{AB}^C$

are not  
covariant

\*) for mathematicians: symplectic reduction of Courant algebroids



# Solution: Poláček-Siegel constr.

produces covariant torsion/curvature  $S$  under gen. Lorentz tr. & gen. diffeomorphisms



# Solution: Poláček-Siegel constr.

produces covariant torsion/curvatureS under gen. Lorentz tr. & gen. diffeomorphisms



2 connections are required:  $\Omega_A^\alpha, \rho^{\alpha\beta}$  ← adjoint index of the gen. Lorentz group  $G_S$

$\Omega_A^\alpha, \rho^{\alpha\beta}, E_A^I$  parameterize a mega-frame  $\mathcal{E}_A^I \in G_{PS}$

# Solution: Poláček-Siegel constr.

produces covariant torsion/curvatureS under gen. Lorentz tr. & gen. diffeomorphisms



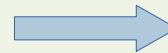
2 connections are required:  $\Omega_A^\alpha, \rho^{\alpha\beta}$  ← adjoint index of the gen. Lorentz group  $G_S$

$\Omega_A^\alpha, \rho^{\alpha\beta}, E_A^I$  parameterize a mega-frame  $\mathcal{E}_A^I \in G_{PS}$

$$O(d+n, d+n) \rightarrow O(d, d) \times G_S, n = \dim(G_S)$$

$$\cup$$

$$G_{PS} \supset G_S$$



# Choosing $G_S$ and $G_{PS}$

## Objective:

1) fix all connections by

1) gauge fixing

2) torsion constraints

in terms of the generalized frame (and its derivatives)

2) as few invariants as possible

We do the same in  
General Relativity.



# Choosing $G_S$ and $G_{PS}$

## Objective:

1) fix all connections by

1) gauge fixing

2) torsion constraints

in terms of the generalized frame (and its derivatives)

2) as few invariants as possible

We do the same in  
General Relativity.

$$G_S = O(d + p) \times O(d + q)$$

$$G_{PS} = O(d + p, d + q)$$

# Recursive embedding of $G_S$

- $G_{PS}$  is generated by  $K_{AB}, R_\alpha^A, R_{\alpha\beta}$
  - and  $G_S$  by  $t_\alpha = (t_{\bar{\alpha}} \leftarrow t_{\underline{\alpha}})$
- How to relate them ???
- left and right factors of  $G_S$

# Recursive embedding of $G_S$

- $G_{PS}$  is generated by  $K_{AB}, R_{\alpha}^A, R_{\alpha\beta}$
  - and  $G_S$  by  $t_{\alpha} = (t_{\bar{\alpha}} \leftarrow t_{\underline{\alpha}})$
- How to relate them ???
- left and right factors of  $G_S$

$$t_{\bar{\alpha}_1} = \frac{1}{\sqrt{a}} (K_{\bar{a}\bar{b}})$$

$$t_{\bar{\alpha}_2} = \frac{1}{2\sqrt{a}} \begin{pmatrix} R_{\bar{\alpha}_1}^{\bar{a}} & R_{\bar{\alpha}_1\bar{\beta}_1} \end{pmatrix}$$

$$t_{\bar{\alpha}_{i+1}} = \frac{1}{2\sqrt{a}} \begin{pmatrix} R_{\bar{\alpha}_i}^{\bar{a}} & R_{\bar{\alpha}_i\bar{\beta}_i} & \dots & R_{\bar{\alpha}_i\bar{\beta}_i} \end{pmatrix}$$

# Recursive embedding of $G_S$

- $G_{PS}$  is generated by  $K_{AB}, R_{\alpha}^A, R_{\alpha\beta}$
  - and  $G_S$  by  $t_{\alpha} = (t_{\bar{\alpha}} \leftarrow t_{\underline{\alpha}})$
- How to relate them ???
- left and right factors of  $G_S$

$$t_{\bar{\alpha}_1} = \frac{1}{\sqrt{a}} (K_{\bar{a}\bar{b}})$$

$$t_{\bar{\alpha}_2} = \frac{1}{2\sqrt{a}} \begin{pmatrix} R_{\bar{\alpha}_1}^{\bar{a}} & R_{\bar{\alpha}_1\bar{\beta}_1} \end{pmatrix}$$

$$t_{\bar{\alpha}_{i+1}} = \frac{1}{2\sqrt{a}} \begin{pmatrix} R_{\bar{\alpha}_i}^{\bar{a}} & R_{\bar{\alpha}_i\bar{\beta}_i} & \dots & R_{\bar{\alpha}_i\bar{\beta}_i} \end{pmatrix}$$

- exponential growth of generators



# Recursive embedding of $G_S$

- $G_{PS}$  is generated by  $K_{AB}, R_{\alpha}^A, R_{\alpha\beta}$
  - and  $G_S$  by  $t_{\alpha} = (t_{\bar{\alpha}} \leftarrow t_{\underline{\alpha}})$
- How to relate them ???
- left and right factors of  $G_S$

$$t_{\bar{\alpha}_1} = \frac{1}{\sqrt{a}} (K_{\bar{a}\bar{b}})$$

$$t_{\bar{\alpha}_2} = \frac{1}{2\sqrt{a}} \begin{pmatrix} R_{\bar{\alpha}_1}^{\bar{a}} & R_{\bar{\alpha}_1\bar{\beta}_1} \end{pmatrix}$$

$$t_{\bar{\alpha}_{i+1}} = \frac{1}{2\sqrt{a}} \begin{pmatrix} R_{\bar{\alpha}_i}^{\bar{a}} & R_{\bar{\alpha}_i\bar{\beta}_i} & \dots & R_{\bar{\alpha}_i\bar{\beta}_i} \end{pmatrix}$$

- exponential growth of generators
- can be truncated at every order

# Torsion constraints and gauge fixing

- Poláček-Siegel construction results one quantity (product):

The twisted mega-space torsion  $\mathcal{T}_{ABC}$  ← fundamental index of  $G_{PS}$

- Like in Cartan geometry, it contains all curvatures and torsions of the gen. connections

$$\Omega_A^\alpha, \rho^{\alpha\beta} .$$

# Torsion constraints and gauge fixing

- Poláček-Siegel construction results one quantity (product):

The twisted mega-space torsion  $\mathcal{T}_{ABC}$  ← fundamental index of  $G_{PS}$

- Like in Cartan geometry, it contains all curvatures and torsions of the gen. connections

$$\Omega_A^\alpha, \rho^{\alpha\beta}.$$

- To fix them completely, we impose:

$$\mathcal{T}_{\underline{ABC}} = \mathcal{T}_{\underline{ABC}} = 0$$

Torsion constraint

$$\Omega_{\underline{a}}^{\bar{\alpha}} = \Omega_{\underline{a}}^{\alpha} = \rho^{\bar{\alpha}\bar{\beta}} = \rho^{\underline{\alpha}\underline{\beta}} = 0$$

Gauge fixing of chiral/anti-chiral sector

# Results

- gauge fixing breaks:

$$G_S \rightarrow \widehat{O}(d) \times \widehat{O}(d)$$

# Results

- gauge fixing breaks:

$$G_S \rightarrow \widehat{O}(d) \times \widehat{O}(d)$$

- after field appropriate field redefinitions, we recover

- 1) deformed Lorentz transformations

(required for Green-Schwarz anomaly cancellation)

- 2) only one invariant

= action (completely check @  $\alpha'$  and partially @  $\alpha'^2$ )

# Results

- gauge fixing breaks:

$$G_S \rightarrow \widehat{O}(d) \times \widehat{O}(d)$$

- after field appropriate field redefinitions, we recover

1) deformed Lorentz transformations

(required for Green-Schwarz anomaly cancellation)

2) only one invariant

= action (completely check @  $\alpha'$  and partially @  $\alpha'^2$ )

There is a hidden symmetry in string theory which controls higher-derivative( $\alpha'$ )-corrections. How far can we push it?